

QE12. Measurement Theory

- Principles of Quantum Measurements

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Lesson #3

In maths...

To define that representation we need a further class of operators: **projection operators** or **projectors** for short. The projector \mathbf{P}_i , onto the eigenstate $|a_i\rangle$ is defined by

$$\mathbf{P}_i := |a_i\rangle\langle a_i|$$

Application of \mathbf{P}_i to an arbitrary state $|\psi\rangle$ yields a multiple of $|a_i\rangle$

$$\mathbf{P}_i|\psi\rangle = |a_i\rangle\langle a_i|\psi\rangle = \langle a_i|\psi\rangle|a_i\rangle$$

where $|\langle a_i|\psi\rangle|$ the “length” of the projection of $|\psi\rangle$ onto the unit vector $|a_i\rangle$. And if $\langle a_i|a_j\rangle = \delta_{ij}$ then

$$\mathbf{P}_i\mathbf{P}_j = \delta_{ij}\mathbf{P}_j; \text{ especially } \mathbf{P}_i^2 = \mathbf{P}_i$$

As the \mathbf{P}_i cover “all directions” of Hilbert space we obtain a completeness relation:

$$\sum_{i=1}^d \mathbf{P}_i = \sum_{i=1}^d |a_i\rangle\langle a_i| = \mathbf{1}$$

Trace of a matrix A : $\text{Tr}(A)$ the sum of the diagonal elements

\mathbf{P}_i is Hermitian: $\mathbf{P}_i = \mathbf{P}_i^\dagger$

Projection postulate

Assume a quantum system prepared in a state $|\psi\rangle$ and a single measurement of the observable \mathbf{A} is performed. This cycle of preparation and measurement is repeated many times so that the notion of probability used in the postulate makes sense.

Or imagine an ensemble containing a large number of independent copies of the quantum system, all prepared in the same state $|\psi\rangle$. \mathbf{A} is measured for all system copies independently.

Projection postulate : A single measurement of the observable \mathbf{A} in the normalized state $|\psi\rangle$ yields one of the eigenvalues a_i of \mathbf{A} with probability $|\langle a_i | \psi \rangle|^2$. Immediately after the measurement the system is in the (normalized) state

$$\frac{\mathbf{P}_i |\psi\rangle}{\|\mathbf{P}_i |\psi\rangle\|}$$

where \mathbf{P}_i is the projection operator onto the subspace of eigenstates of \mathbf{A} with eigenvalue a_i .

Projective measurement result

In general it is not possible to predict the outcome of a single measurement. A measurement of \mathbf{A} on an ensemble of systems as discussed above yields the *average* (expectation value)

$$\langle \mathbf{A} \rangle := \langle \psi | \mathbf{A} | \psi \rangle$$

with deviations described by the **variance** (the square of the standard deviation): $\langle (\mathbf{A} - \langle \mathbf{A} \rangle)^2 \rangle \geq 0$

The probability of obtaining outcome i for a given state $|\psi\rangle$

$$p_i = \langle \psi | \mathbf{P}_i | \psi \rangle$$

And the post-measurement state is given by

$$|\psi_i^{post}\rangle = \frac{\mathbf{P}_i |\psi\rangle}{\sqrt{\langle \psi | \mathbf{P}_i | \psi \rangle}}$$

In quantum mechanics the measurement change the state of a quantum system which is **probabilistic** and **irreversible** process. **The observation process is irreversible.**

Projectors: a methodology for a single qubit

1. Identify the projectors for the basis:

$$P_i = |v_i\rangle\langle v_i|$$

2. Compute each projected vector:

$$|\psi_{P_i}\rangle = P_i |\psi\rangle$$

3. Compute the squared magnitude of each projection: as the probability of seeing the i-th normalized projection.

$$\|\psi_{P_i}\|^2$$

4. Compute each normalized projection (each potential outcome):

$$|\psi_{N_i}\rangle = \frac{|\psi_{P_i}\rangle}{\|\psi_{P_i}\|}$$

This gives us each possible outcome and its associated probability:

Outcome $|\psi_{N_i}\rangle$ occurs with probability $\|\psi_{P_i}\|^2$

Example 1:

Suppose we measure the state $|0\rangle$ in the basis $\{|+\rangle, |-\rangle\}$.

$$|0\rangle = \frac{1}{\sqrt{2}} (|+\rangle + |-\rangle)$$

Step 1: Identify the projectors for the basis.

$$P_+ = |+\rangle\langle+|$$

$$P_- = |-\rangle\langle-|$$

Step 2: Compute each projected vector:

$$P_+ |0\rangle = |+\rangle\langle+|0\rangle = |+\rangle\langle+| \left(\frac{1}{\sqrt{2}}|+\rangle + \frac{1}{\sqrt{2}}|-\rangle \right) = \frac{1}{\sqrt{2}}|+\rangle$$

$$P_- |0\rangle = |-\rangle\langle-|0\rangle = |-\rangle\langle-| \left(\frac{1}{\sqrt{2}}|+\rangle + \frac{1}{\sqrt{2}}|-\rangle \right) = \frac{1}{\sqrt{2}}|-\rangle$$

Example 1:

Step 3: Compute the squared magnitude of each projection:

$$\begin{aligned} |P_+ |0\rangle|^2 &= \left| \frac{1}{\sqrt{2}} |+\rangle \right|^2 = \left\langle \frac{1}{\sqrt{2}} \langle + | \left| \frac{1}{\sqrt{2}} |+\rangle \right\rangle = \frac{1}{2} \\ |P_- |0\rangle|^2 &= \left| \frac{1}{\sqrt{2}} |-\rangle \right|^2 = \left\langle \frac{1}{\sqrt{2}} \langle - | \left| \frac{1}{\sqrt{2}} |-\rangle \right\rangle = \frac{1}{2} \end{aligned}$$

Step 4: Compute each normalized projection (each potential outcome):

$$\frac{\frac{1}{\sqrt{2}} |+\rangle}{\left| \frac{1}{\sqrt{2}} |+\rangle \right|} = |+\rangle \quad \frac{\frac{1}{\sqrt{2}} |-\rangle}{\left| \frac{1}{\sqrt{2}} |-\rangle \right|} = |-\rangle$$

This gives us possible outcomes and their associated probabilities:

outcome $|+\rangle$ occurs with probability $\frac{1}{2}$
outcome $|-\rangle$ occurs with probability $\frac{1}{2}$

Example: two-qubits states

Consider the two-qubit state $|\psi\rangle = \frac{1}{\sqrt{2}}|01\rangle + \frac{1}{\sqrt{2}}|10\rangle$

And suppose we use the two-qubit standard basis for two-qubit measurement: $|00\rangle$, $|01\rangle$, $|10\rangle$, $|11\rangle$

Step 1: Identify the measurement basis projectors:

$$\begin{aligned} P_{00} &= |00\rangle\langle 00| \\ P_{01} &= |01\rangle\langle 01| \\ P_{10} &= |10\rangle\langle 10| \\ P_{11} &= |11\rangle\langle 11| \end{aligned}$$

For 2 qubits in the standard basis, the shorthand is:

$$\begin{aligned} |0\rangle \otimes |0\rangle &= |00\rangle \\ |0\rangle \otimes |1\rangle &= |01\rangle \\ |1\rangle \otimes |0\rangle &= |10\rangle \\ |1\rangle \otimes |1\rangle &= |11\rangle \end{aligned}$$

Example: two-qubits states

Step 2: Use the above to compute projected vectors:

$$\begin{aligned}P_{00} |\psi\rangle &= |00\rangle\langle 00| \left(\frac{1}{\sqrt{2}}|01\rangle + \frac{1}{\sqrt{2}}|10\rangle \right) = 0 \\P_{01} |\psi\rangle &= |01\rangle\langle 01| \left(\frac{1}{\sqrt{2}}|01\rangle + \frac{1}{\sqrt{2}}|10\rangle \right) = \frac{1}{\sqrt{2}}|01\rangle \\P_{10} |\psi\rangle &= |10\rangle\langle 10| \left(\frac{1}{\sqrt{2}}|01\rangle + \frac{1}{\sqrt{2}}|10\rangle \right) = \frac{1}{\sqrt{2}}|10\rangle \\P_{11} |\psi\rangle &= |11\rangle\langle 11| \left(\frac{1}{\sqrt{2}}|01\rangle + \frac{1}{\sqrt{2}}|10\rangle \right) = 0\end{aligned}$$

Step 3: The magnitudes of (non-zero) projected vectors are their probabilities:

$$\begin{aligned}|P_{01} |\psi\rangle|^2 &= \left| \frac{1}{\sqrt{2}}|01\rangle \right|^2 = \frac{1}{2} \\|P_{10} |\psi\rangle|^2 &= \left| \frac{1}{\sqrt{2}}|10\rangle \right|^2 = \frac{1}{2}\end{aligned}$$

Step 4: Normalize the projections to get the outcome vectors:

$$\begin{aligned}\frac{1}{|P_{01} |\psi\rangle|} P_{01} |\psi\rangle &= |01\rangle \\ \frac{1}{|P_{10} |\psi\rangle|} P_{10} |\psi\rangle &= |10\rangle\end{aligned}$$

Tip for calculations

Notice the usefulness of Dirac notation, over matrices, in the second step, for example:

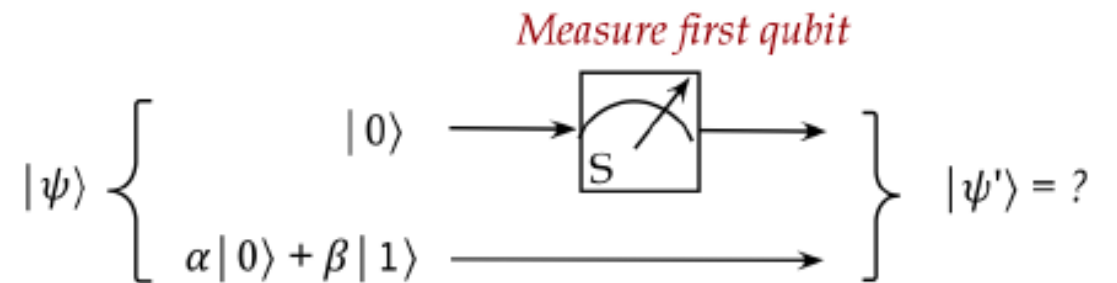
$$\begin{aligned} & |01\rangle\langle 01| \left(\frac{1}{\sqrt{2}} |01\rangle + \frac{1}{\sqrt{2}} |10\rangle \right) \\ &= \left(\frac{1}{\sqrt{2}} |01\rangle \langle 01| |01\rangle + \frac{1}{\sqrt{2}} |01\rangle \langle 01| |10\rangle \right) \\ &= \frac{1}{\sqrt{2}} |01\rangle \quad \text{inner-product} = 1 \quad \text{inner-product} = 0 \end{aligned}$$

Scalar movement (indicated by a red arrow pointing to the scalar $\frac{1}{\sqrt{2}}$ in the second line)

Example #3

Suppose you have a 2-qubit state and we want to measure only the first qubit. The input two-qubit vector is $|\psi\rangle = |0\rangle \otimes (\alpha|0\rangle + \beta|1\rangle)$

After measuring only the first qubit, what are the possible states, and with what probabilities?



Solution:

Clearly, since the first (top) qubit is already $|0\rangle$, measuring this qubit alone should leave the top qubit as $|0\rangle$. Intuition suggests that because there's no entanglement, the second qubit should be the same. That is, after measurement, that state should be

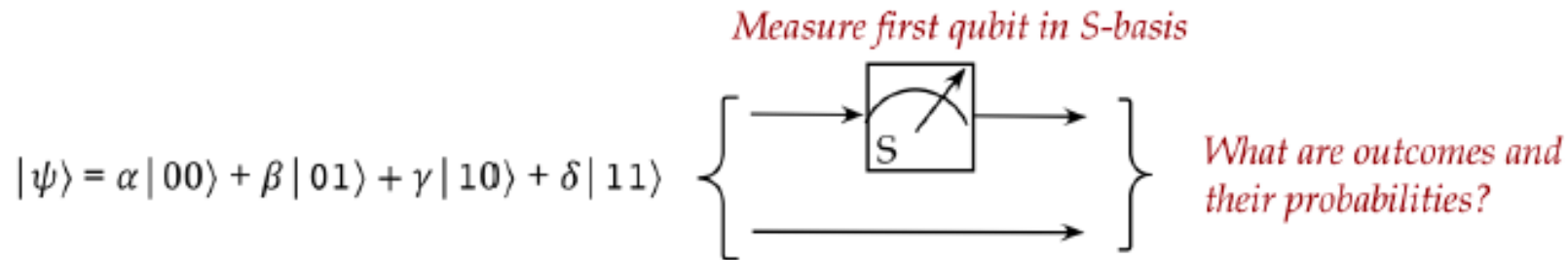
$$|\psi'\rangle = |0\rangle \otimes (\alpha|0\rangle + \beta|1\rangle) \quad \text{Or} \quad |\psi'\rangle = \alpha(|0\rangle \otimes |0\rangle) + \beta(|0\rangle \otimes |1\rangle) = \alpha|00\rangle + \beta|01\rangle$$

Projective measurement: extending the theory to multiple qubits

We want to extend the theory to be able to handle different measurement scenarios with qubits.

Examples of scenarios include:

- Measuring all qubits simultaneously.
- Measuring just one qubit amongst the qubits.
- Measuring a subset of qubits from the qubits.
- And each case, having the freedom to use a variety of measurement bases.
- There are three aspects to extending single-qubit projective measurement to multiple qubits:
- Understanding how the particular measurement splits the whole n -qubit space into orthogonal subspaces.
- Building n -qubit projectors accordingly.
- Seeing if it helps to construct the n -qubit projectors from smaller projectors (such as 1-qubit projectors).



$$|\psi\rangle = \alpha|00\rangle + \beta|01\rangle + \gamma|10\rangle + \delta|11\rangle$$

We want to measure the one qubit.

We will go about addressing this question in several stages:

- We'll first examine the potential vectors that result from subjecting the 2-qubit vector to 1-qubit measurement.
- These potential vectors will form a vector space.
- Then, we'll build the projectors for these spaces.
- We'll also see that the projectors can be built from tensoring.

Stage 1: if you measure the first qubit, what are the possible 2-qubit vectors that result?

Example: first qubit results in $|0\rangle$.

Second qubit could potentially be any state.

Describe this as the space

$$V_1 = \text{span} \{ |00\rangle, |01\rangle \}$$

Any vector in this space has the first qubit as $|0\rangle$.

Similarly, the space of 2-qubit vectors that correspond to "first qubit is $|1\rangle$ " is

$$V_2 = \text{span} \{ |10\rangle, |11\rangle \}$$

Stage 2: Thus, the potential results of first-qubit measurement lie in two orthogonal subspaces:

$$V = V_1 \cup V_2$$

where

$$V_1 = \text{span} \{ |00\rangle, |01\rangle \} \quad \text{First qubit } |0\rangle$$

$$V_2 = \text{span} \{ |10\rangle, |11\rangle \} \quad \text{First qubit } |1\rangle$$

Stage 3: Now define projectors for each of these subspaces

$$P_{V_1} = |00\rangle\langle 00| + |01\rangle\langle 01|$$

$$P_{V_2} = |10\rangle\langle 10| + |11\rangle\langle 11|$$

Note: applying in Dirac form makes it easy to see why these are projectors for those subspaces:

- Consider any vector

$$|\psi\rangle = \alpha|00\rangle + \beta|01\rangle + \gamma|10\rangle + \delta|11\rangle$$

- Then

$$\begin{aligned} P_{V_1} |\psi\rangle &= \text{projection of } |\psi\rangle \text{ on } V_1 \\ &= (|00\rangle\langle 00| + |01\rangle\langle 01|) (\alpha|00\rangle + \beta|01\rangle + \gamma|10\rangle + \delta|11\rangle) \\ &= \alpha|00\rangle + \beta|01\rangle \\ &= \text{A vector in } V_1 \end{aligned}$$

Recall: $V_1 = \text{span}\{|00\rangle, |01\rangle\}$

- Similarly,

$$\begin{aligned} P_{V_2} |\psi\rangle &= \text{projection of } |\psi\rangle \text{ on } V_2 \\ &= \gamma|10\rangle + \delta|11\rangle \\ &= \text{A vector in } V_2 \end{aligned}$$

The *outcomes* of measurement and their probabilities are

$$\begin{aligned} \text{normalized } P_{V_1} |\psi\rangle &\text{ occurs with probability } |P_{V_1} |\psi\rangle|^2 \\ \text{normalized } P_{V_2} |\psi\rangle &\text{ occurs with probability } |P_{V_2} |\psi\rangle|^2 \end{aligned}$$

- Squared magnitudes of the projections are:

$$\begin{aligned} |P_{V_1} |\psi\rangle|^2 &= |\alpha|^2 + |\beta|^2 \\ |P_{V_2} |\psi\rangle|^2 &= |\gamma|^2 + |\delta|^2 \end{aligned}$$

- Normalized projections are:

$$\begin{aligned} \frac{1}{|P_{V_1} |\psi\rangle|} P_{V_1} |\psi\rangle &= \frac{1}{\sqrt{|\alpha|^2 + |\beta|^2}} (\alpha |00\rangle + \beta |01\rangle) \\ \frac{1}{|P_{V_2} |\psi\rangle|} P_{V_2} |\psi\rangle &= \frac{1}{\sqrt{|\gamma|^2 + |\delta|^2}} (\gamma |10\rangle + \delta |11\rangle) \end{aligned}$$



Let's take a closer look at the first projected vector to see that it makes sense:

- Recall we started in state

$$|\psi\rangle = \alpha|00\rangle + \beta|01\rangle + \gamma|10\rangle + \delta|11\rangle$$

- Let's write this as

$$|\psi\rangle = |0\rangle \otimes (\alpha|0\rangle + \beta|1\rangle) + |1\rangle \otimes (\gamma|0\rangle + \delta|1\rangle)$$

Here, we've just separated out the first qubit for emphasis.

- If we measure the first qubit as $|0\rangle$, then the second should be "untouched" in state $\alpha|0\rangle + \beta|1\rangle$.
- And the resulting 2-qubit vector should be

$$|0\rangle \otimes (\alpha|0\rangle + \beta|1\rangle) = \alpha|00\rangle + \beta|01\rangle$$

- But the latter is not normalized, and so, the normalized vector is:

$$\frac{1}{\sqrt{|\alpha|^2 + |\beta|^2}}(\alpha|00\rangle + \beta|01\rangle)$$

Can the 2-qubit projectors be built out of smaller 1-qubit projectors?

Recall that the two 1-qubit S-basis projectors are

$$\begin{aligned}P_0 &= |0\rangle\langle 0| \\P_1 &= |1\rangle\langle 1|\end{aligned}$$

Since we're not measuring the 2nd-qubit, the only projector that keeps the qubit the same is the identity

$$I = |0\rangle\langle 0| + |1\rangle\langle 1|$$

(in Dirac form).

Thus, one can construct the 2-qubit projectors via tensoring

$$\begin{aligned}P_{V_1} &= P_0 \otimes I = |0\rangle\langle 0| \otimes (|0\rangle\langle 0| + |1\rangle\langle 1|) \\P_{V_2} &= P_1 \otimes I = |1\rangle\langle 1| \otimes (|0\rangle\langle 0| + |1\rangle\langle 1|)\end{aligned}$$

Let's work out the first one to see the details:

$$\begin{aligned} P_0 \otimes I &= |0\rangle\langle 0| \otimes (|0\rangle\langle 0| + |1\rangle\langle 1|) \\ &= (|0\rangle\langle 0| \otimes |0\rangle\langle 0|) + (|0\rangle\langle 0| \otimes |1\rangle\langle 1|) \\ &= (|0 \otimes 0\rangle\langle 0 \otimes 0|) + (|0 \otimes 1\rangle\langle 0 \otimes 1|) \\ &= |00\rangle\langle 00| + |01\rangle\langle 01| \end{aligned}$$

Tensor of projectors

Tensor properties

Proposition 4.5

Shorthand notation

Please prove that:

$$|v\rangle\langle v| \otimes |w\rangle\langle w| = |v \otimes w\rangle\langle v \otimes w|$$

Finally, let's remind ourselves of alternative tensoring notation:

- We can write

$$|v\rangle\langle v| \otimes |w\rangle\langle w| = |\mathbf{v}\rangle|\mathbf{w}\rangle\langle\mathbf{v}|\langle\mathbf{w}| \quad (\text{Also written as } |v, w\rangle\langle v, w|)$$

- Thus, we could also have written the earlier projector example as:

$$\begin{aligned} P_0 \otimes I &= |0\rangle\langle 0| \otimes (|0\rangle\langle 0| + |1\rangle\langle 1|) \\ &= (|0\rangle\langle 0| \otimes |0\rangle\langle 0|) + (|0\rangle\langle 0| \otimes |1\rangle\langle 1|) \\ &= (|\mathbf{0}\rangle|\mathbf{0}\rangle\langle\mathbf{0}|\langle\mathbf{0}|) + (|\mathbf{0}\rangle|\mathbf{1}\rangle\langle\mathbf{0}|\langle\mathbf{1}|) \\ &= |00\rangle\langle 00| + |01\rangle\langle 01| \end{aligned}$$

- By convention, we do *not* write $|00\rangle$ as $|0, 0\rangle$.

Homework

Consider the single-qubit projectors $P_+ = |+\rangle\langle+|$ and $P_- = |-\rangle\langle-|$.

1. Show that $(P_+ \otimes I)|00\rangle = \frac{1}{\sqrt{2}}|+, 0\rangle$ and $(P_+ \otimes I)|11\rangle = \frac{1}{\sqrt{2}}|+, 1\rangle$
2. Show that $\frac{1}{\sqrt{2}}|+, 0\rangle + \frac{1}{\sqrt{2}}|+, 1\rangle = |+, +\rangle$
3. Expand $(P_+ \otimes I)$ into a matrix and use that to show $(P_+ \otimes I)|\psi\rangle = |+, +\rangle$ (normalized) where $|\psi\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$

End of Lesson